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1999 J. Phys. A: Math. Gen. 32 443

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ADDENDUM

Duality for multiparametric quantum GL(n)

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Received 12 October 1998

Abstract. In our paper (Dobrev V K and Parashar P 1993 *J. Phys. A: Math. Gen.* **26** 6991–7002) we have constructed the Hopf algebra \mathcal{U} as dual to the quantum group deformation of GL(n) (depending on 1 + n(n-1)/2 parameters), and we have shown that \mathcal{U} has one-generator central Hopf subalgebra and a commutation subalgebra \mathcal{U}' which is a Drinfeld–Jimbo-like deformation of U(sl(n)), but which is not a Hopf subalgebra since its coproducts depend also on the central generator. In this addendum we show that \mathcal{U}' is a Hopf subalgebra, besides the trivial case of coinciding parameters, also for a special choice of the parameters, so that there would be only 1 + (n-1)(n-2)/2 independent parameters.

In the main text [1] we have constructed the Hopf algebra $\mathcal{U} \equiv \mathcal{U}_{uq}$ as the dual to the multiparameter matrix quantum group $GL_{uq}(n)$ (depending on the parameter u and n(n-1)/2 parameters $q = \{q_{ij} | j - i > 0\}$) and we have shown that \mathcal{U} may be split in the form: $U_{u,q}(sl(n, \mathbb{C})) \otimes U_u(\mathcal{Z})$, where \mathcal{Z} is one-dimensional spanned by the generator K (cf (29)) and $U_u(\mathcal{Z})$ is a central Hopf subalgebra of \mathcal{U} , while $\mathcal{U}' = U_{u,q}(sl(n, \mathbb{C}))$ is a deformation of $U(sl(n, \mathbb{C}))$ which is of Drinfeld–Jimbo form with deformation parameter u as a commutation subalgebra of \mathcal{U} . However, \mathcal{U}' is not a Hopf subalgebra of \mathcal{U} since the coproducts of \mathcal{U}' depend also on the generator K. Furthermore, these coproducts depend on all parameters u and q. We have also noticed that \mathcal{U}' is a Hopf subalgebra if all parameters coincide.

In this addendum we show that \mathcal{U}' is a Hopf subalgebra also for a special choice of the parameters, so that there would be only 1 + (n - 1)(n - 2)/2 independent parameters: u and $\tilde{q} = \{q_{ij} | j - i > 1\}$. In this special case the central generator K decouples from the coproducts of the generators of $\mathcal{U}' = U_{u,\tilde{q}}(sl(n, \mathbb{C}))$ as in the one-parameter deformation, however, these coproducts depend also on \tilde{q} .

To make this explicit we first need to express the operators \mathcal{P}_i (cf (43), and through them \mathcal{Q}_i , cf (45)) in terms of the generators H_i and K. For this we first express the generators D_i through H_i and K:

$$D_{i} = \frac{1}{n} \left(K - \sum_{j=1}^{i-1} j H_{j} + \sum_{j=i}^{n-1} (n-j) H_{j} \right) = \hat{K} + \hat{H}_{i}$$

$$\hat{K} \equiv \frac{1}{n} \left(K - \sum_{j=1}^{n-1} j H_{j} \right) \qquad \hat{H}_{i} \equiv \sum_{j=i}^{n-1} H_{j} \qquad (\hat{H}_{n} \equiv 0).$$
(A.1)

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0305-4470/99/020443+02\$19.50 © 1999 IOP Publishing Ltd

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Now we substitute (A.1) into (43) to obtain:

$$\mathcal{P}_{i} = (\tilde{q}_{i})^{\hat{K}} \left(\prod_{s=1}^{i-1} \left(\frac{q_{si}}{q_{s,i+1}} \right)^{\hat{H}_{s}} \right) \left(\frac{u^{2}}{q_{i,i+1}} \right)^{\hat{H}_{i}} \left(\frac{1}{q_{i,i+1}} \right)^{\hat{H}_{i+1}} \prod_{t=i+2}^{n-1} \left(\frac{q_{i+1,t}}{q_{it}} \right)^{\hat{H}_{t}}$$

$$\tilde{q}_{i} \equiv \left(\prod_{s=1}^{i-1} \frac{q_{si}}{q_{s,i+1}} \right) \frac{u^{2}}{q_{i,i+1}^{2}} \prod_{t=i+2}^{n} \frac{q_{i+1,t}}{q_{it}}.$$
(A.2)

From the above expression it is clear that in order for K to decouple from the system the n-1 constants \tilde{q}_i should become equal to unity. This brings n-1 conditions on the parameters q_{ij} . It seems natural to use these conditions to fix the n-1 next-to-main-diagonal parameters $q_{i,i+1}$, and indeed, a natural choice for this exists, namely, we may set:

$$q_{i,i+1}^{0} \equiv u^{i(n-i)} \prod_{\substack{1 \le s \le i \\ s < t-1}} q_{st}^{-1}$$

= $u \prod_{s=1}^{i} \prod_{t=i+1}^{n} \frac{u}{q_{st}}$ $1 \le i \le n-1$ (A.3)

where the tilde over the double product means that the case s = i = t - 1 should be omitted. Then we obtain:

$$(\tilde{q}_i)_{q_{i,i+1}=q_{i,i+1}^0} = 1$$
 $1 \le i \le n-1$ (A.4)

and substituting this in the operators \mathcal{P}_i we get in terms of \hat{H}_i and in terms of H_i :

$$\tilde{\mathcal{P}}_{i} = (\mathcal{P}_{i})_{q_{i,i+1}=q_{i,i+1}^{0}} = \left(\prod_{s=1}^{i-2} \left(\frac{q_{si}}{q_{s,i+1}}\right)^{\hat{H}_{s}}\right) \left(\frac{u}{q_{i-1,i+1}} \prod_{s=1}^{i-1} \prod_{t=i}^{n} \frac{u}{q_{st}}\right)^{\hat{H}_{i-1}} \times u^{\hat{H}_{i}-\hat{H}_{i+1}} \left(\prod_{s=1}^{i} \prod_{t=i+1}^{n} \frac{q_{st}}{u}\right)^{\hat{H}_{i}+\hat{H}_{i+1}} \left(\frac{u}{q_{i,i+2}} \prod_{s=1}^{i+1} \prod_{t=i+2}^{n} \frac{u}{q_{st}}\right)^{\hat{H}_{i+2}} \prod_{t=i+3}^{n-1} \left(\frac{q_{i+1,t}}{q_{it}}\right)^{\hat{H}_{t}} = \left(\prod_{j=1}^{i-2} \left(\prod_{s=1}^{j} \frac{q_{si}}{q_{s,i+1}}\right)^{H_{j}}\right) \left(\left(\prod_{s=1}^{i-1} \frac{u^{2}}{q_{s,i+1}^{2}}\right) \prod_{s=1}^{i-1} \prod_{t=i+2}^{n} \frac{u}{q_{st}}\right)^{H_{i-1}} \times \left(u \left(\prod_{s=1}^{i-1} \frac{u}{q_{s,i+1}}\right) \prod_{t=i+2}^{n} \frac{q_{it}}{u}\right)^{H_{i}} \left(\left(\prod_{t=i+2}^{n} \frac{q_{it}^{2}}{u^{2}}\right) \prod_{s=1}^{i-1} \prod_{t=i+2}^{n} \frac{q_{st}}{u}\right)^{H_{i+1}} \times \left(\prod_{j=i+2}^{n-1} \left(\prod_{t=j+1}^{n} \frac{q_{it}}{q_{i+1,t}}\right)^{H_{j}}\right). \tag{A.5}$$

Thus, for the particular choice $q_{i,i+1} = q_{i,i+1}^0$ we have the splitting $\mathcal{U}_{u\tilde{q}} \cong U_{u,\tilde{q}}(sl(n, \mathbb{C})) \otimes U_u(\mathcal{Z})$ as tensor product of two Hopf subalgebras. Note that $U_{u,\tilde{q}}(sl(n, \mathbb{C}))$ is a Hopf algebra which is a deformation of $U(sl(n, \mathbb{C}))$, and is of Drinfeld–Jimbo form with deformation parameter *u* as commutation algebra, but not as a coalgebra since the coproducts depend also on \tilde{q} .

References

[1] Dobrev V K and Parashar P 1993 J. Phys. A: Math. Gen. 26 6991-7002